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CONTACT HAMILTONIANS DISTINGUISHING LOCALLY CERTAIN GOURSAT SYSTEMS

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Abstract. For the first time in dimension 9, the Goursat distributions are not locally smoothly classified by their small growth vector at a point. As shown in [M1], in dimension 9 of the underlying manifold 93 different local behaviours are possible and four irregular pairs of them have coinciding small growth vectors. In the present paper we distinguish geometrically objects in three of those pairs. Smooth functions in three variables – contact hamiltonians in the terminology of Arnold, [A] – help to do that. One pair of models, however, resists this technique. Another example of similar resistance in dimension 10 is also given – through the exact classification in dimension 10 of one family of local pseudo-normal forms (with redundant real constants) for Goursat objects. The latter result is an harbinger of more general phenomena that will be treated in a subsequent paper.

1. Introduction. For any geometric distribution D (a subbundle of the tangent bundle) on a given smooth manifold M we set $D_1 = D$, $D_2 = D + [D, D], \ldots, D_{l+1} = D_l + [D, D_l];$ $D^{(0)} = D$, $D^{(1)} = D + [D, D], \ldots, D^{(l+1)} = D^{(l)} + [D^{(l)}, D^{(l)}]$. All these are, in general, modules of vector fields. By the *small growth vector* of D at p we understand the sequence $[n_1, n_2, n_3, \ldots]$ of dimensions at $p \in M$ of the flag $D_1 \subset D_2 \subset D_3 \subset \ldots$

DEFINITION 1. Let D be a rank-2 distribution on an n-dimensional manifold. We say that D satisfies the *Goursat Condition* (GC for short in the sequel) when the members of the flag $D^{(0)} \subset D^{(1)} \subset D^{(2)} \subset \ldots$ have at every point linear dimensions 2, 3, ..., n-1, n. The dual object $S = D^{\perp}$ we just call a *Goursat system*.

This condition is sometimes also called the Cartan–Goursat condition.

Every Goursat system S on an n-dimensional manifold admits locally around any fixed point certain Kumpera-Ruiz pseudo-normal form (see [KR] and, for newer presentations,

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[CM], [PR]), depending on that point and not unique (in general) even after fixing the point: the germ at $0 \in \mathbf{R}^n(x^1, x^2, \ldots, x^n)$ of $(\omega^1, \omega^2, \ldots, \omega^{n-2})$, where

$$\begin{split} \omega^{1} &= dx^{i_{1}} - x^{3}dx^{j_{1}}, \quad (i_{1}, j_{1}) = (2, 1) \\ \omega^{2} &= dx^{i_{2}} - x^{4}dx^{j_{2}}, \quad (i_{2}, j_{2}) = (3, j_{1}) \\ \omega^{3} &= dx^{i_{3}} - x^{5}dx^{j_{3}}, \quad (i_{3}, j_{3}) \in \{(4, j_{2}), (j_{2}, 4)\} \\ \omega^{4} &= dx^{i_{4}} - X^{6}dx^{j_{4}}, \quad (i_{4}, j_{4}) \in \{(5, j_{3}), (j_{3}, 5)\} \\ \bullet & \bullet \\ \omega^{n-2} &= dx^{i_{n-2}} - X^{n}dx^{j_{n-2}}, (i_{n-2}, j_{n-2}) \in \{(n-1, j_{n-3}), (j_{n-3}, n-1)\}, \end{split}$$

and where for $6 \le l \le n$, $X^{l} = x^{l}$ if $(i_{l-2}, j_{l-2}) = (j_{l-3}, l-1)$ and $X^{l} = x^{l} + c^{l}$ in the opposite case of $(i_{l-2}, j_{l-2}) = (l-1, j_{l-3})$. The $c^{6}, c^{7}, \ldots, c^{n}$ are certain real constants.¹

Moreover, one gets locally the successive derived systems of S by always removing the last (bottommost) Pfaffian equation.

(Throughout this paper, in contrast to [KR] and [CM], we write the forms ω^l with the – signs instead of +. The same convention was used in [G].)

We code the KR pseudo-normal forms in dimensions $n \ge 5$ as in [CM], assigning to a Pfaffian equation $\omega^l = 0, l = 3, 4, ..., n - 2$:

- 1, when the first alternative for ω^l holds and $c^{l+2} = 0$,
- 2, in the case of the first alternative, but $c^{l+2} \neq 0$,
- 3, in the case of the second alternative for ω^l ,

and writing those integers, when l runs through the set $\{3, 4, \ldots, n-2\}$, from left to right, separated by dots. Specifically, we write **2** instead of 2 when the relevant constant is 1, and **2**- when it is -1.

EXAMPLE 1. The germ at $0 \in \mathbf{R}^9(x^1, x^2, \dots, x^9)$ of $(\omega^1, \omega^2, \dots, \omega^7)$, where $\omega^1 = dx^2 - x^3 dx^1$, $\omega^2 = dx^3 - x^4 dx^1$, $\omega^3 = dx^1 - x^5 dx^4$, $\omega^4 = dx^5 - (1+x^6)dx^4$, $\omega^5 = dx^4 - x^7 dx^6$, $\omega^6 = dx^6 - x^8 dx^7$, $\omega^7 = dx^8 - (c+x^9)dx^7$ with c = 1 (-1) is coded 3.2.3.3.2 (3.2.3.3.2-, respectively). These two local models form the exceptional pair (***) listed in [M1], Main Theorem.

We are going to consider KR pseudo-normal forms in arbitrary dimension $n \geq 3$ not as germs at 0 but as representatives defined on the whole \mathbf{R}^n . By an *automorphism* of such a differential system S we mean any diffeomorphism $g: \mathbf{R}^n \leftrightarrow$ preserving $S: g^*S = S$. By an *infinitesimal automorphism* (i. a.) of S we mean any smooth vector field on \mathbf{R}^n whose flow (at least for small |t|) preserves S. The set of all i. a.'s of S will be denoted – following the tradition established in [KR], [G] and [K] – by L(S).

2. Contact hamiltonians parametrizing i. a.'s of KR pseudo-normal forms

2.1. Dimension 3. What are i.a.'s of the Darboux contact structure $\omega^1 = dx^2 - x^3 dx^1 = 0$ on $\mathbf{R}^3(x^1, x^2, x^3)$? The answer was given already by S. Lie [Li]. To reproduce

¹The 'pseudo' refers to the fact that the constants are not, in general, invariants of the local classification.

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it here, following the exposition in [Ly] (done for contact structures in arbitrary odd dimension, in fact), let us replace the coordinates x^1 , x^2 , x^3 by the coordinates q^1 , u, p^1 in the bundle $J^1(\mathbf{R}^1)$ of the 1-jets of scalar functions $u(q^1)$ of one variable q^1 , with $p^1 = \frac{du}{dq^1}$. The contact distribution $(dx^2 - x^3 dx^1)^{\perp}$ becomes then the *Cartan distribution* \mathcal{C} on $J^1(\mathbf{R}^1)^2$ deciding which sections of this bundle come from genuine functions of one variable (those that are tangent to \mathcal{C}). All i. a.'s of \mathcal{C} are of the form

$$-\frac{\partial f}{\partial p^1}\frac{\partial}{\partial q^1} + \left(f - p^1\frac{\partial f}{\partial p^1}\right)\frac{\partial}{\partial u} + \left(\frac{\partial f}{\partial q^1} + p^1\frac{\partial f}{\partial u}\right)\frac{\partial}{\partial p^1}$$

where $f = f(q^1, u, p^1)$ is an arbitrary smooth function. Returning to our notation and writing ∂_k for $\partial/\partial x^k$, all i. a.'s of (\mathbf{R}^3, ω^1) are of the form

$$-f_3 \partial_1 + (f - x^3 f_3) \partial_2 + (f_1 + x^3 f_2) \partial_3.$$
 (1)

An arbitrary smooth function $f = f(x^1, x^2, x^3)$ in (1) is called a *contact hamiltonian* because of a natural (and deep) interpretation in terms of the symplectization of any given contact manifold – see [A], App. 4.

Expression (1) for i. a.'s of Darboux is used, omitting references, in [G] and [K].

2.2. Parametrization of dimension $n \ge 4$ by dimension 3. We recall in this section that when a KR pseudo-normal form S is given on $\mathbf{R}^n(x^1, \ldots, x^n)$, its first derived system $S^{(1)}$ is defined on $\mathbf{R}^{n-1}(x^1, \ldots, x^{n-1})$ and suspended only in the direction ∂_n , according to the information quoted in Sec. 1. In consequence, abusing notation a bit, members of $L(S^{(1)})$ are understood as vector fields on \mathbf{R}^{n-1} only.

PROPOSITION 1 ([G], [K]). (i) Every $X \in L(S)$ has all but the last coordinate functions depending only on x^1, \ldots, x^{n-1} . In this way, denoting by $[X]_1$ the truncation of X by its last coordinate function, we have $[X]_1 \in L(S^{(1)})$.

(ii) Every $Y \in L(S^{(1)})$ has a unique prolongation $[Y]^1 \in L(S)$ s.t. $[[Y]^1]_1 = Y$. (iii) In the situation of (i), by the uniqueness in (ii), $[[X]_1]^1 = X$.

PROOF. (i). The key factor is Cor. 3.1 of [M1].³ Because of that, the flow φ_X^t of X can be written as

$$\varphi_X^t(x^1, x^2, \dots, x^n) = (f^1(t, x^1, \dots, x^{n-1}), \dots, f^{n-1}(t, x^1, \dots, x^{n-1}), f^n(t, x^1, \dots, x^n)).$$

Therefore, writing horizontally the coordinates of a vector field,

$$X(x^{1}, x^{2}, \dots, x^{n}) = \left[\frac{d}{dt}f^{1}|_{t=0}, \dots, \frac{d}{dt}f^{n-1}|_{t=0}, \frac{d}{dt}f^{n}|_{t=0}\right].$$

This shows that all but last coordinate functions of X depend only on t, x^1, \ldots, x^{n-1} , hence only on x^1, \ldots, x^{n-1} , X being time-independent.

(ii). We seek prolongation(s) of $Y, Y + F\partial_n$ preserving (infinitesimally) the new Pfaffian equation $\omega^{n-2} = 0$ modulo $(\omega^1, \ldots, \omega^{n-3})$:

$$L_{Y+F\partial_n}(\omega^{n-2}) \in (\omega^1, \dots, \omega^{n-2}).$$
⁽²⁾

²Lychagin calls $du - \sum_{i=1}^{k} p^{i} dq^{i}$ a classifying element on $J^{1}(\mathbf{R}^{k})$.

³Stated there for diffeos g preserving 0, but with the condition g(0) = 0 not used in the proof. That fact can be, essentially, retraced already in E. Cartan's work [C], p. 10.

Because the new variable (in the form of x^n or $X^n = c^n + x^n$) enters the new equation, while its derivative dx^n does not, equation (2) is always a linear equation for F with leading coefficient 1. A careful inspection of the KR pseudo-normal forms shows that the free term uses just the first derivatives of coordinate functions of Y. Instead of abstract writing, let us analyze in detail the initial cases n = 4 and 5. The general rules of prolongation do not differ much from them; as a consequence, the existence and uniqueness of F follow.

n = 4; we write $Y = A\partial_1 + B\partial_2 + C\partial_3$, and name $D\partial_4$ the term prolonging Y. Then (2) reads

$$dC - x^4 dA - D dx^1 = (C_2 - x^4 A_2)(dx^2 - x^3 dx^1) + (C_3 - x^4 A_3)(dx^3 - x^4 dx^1),$$

implying $D = C_1 - x^4 A_1 + x^3 (C_2 - x^4 A_2) + x^4 (C_3 - x^4 A_3).$

n = 5; $Y = A\partial_1 + B\partial_2 + C\partial_3 + D\partial_4$, the prolonging term denoted now by $E\partial_5$. If S is 1, i. e., $\omega^3 = dx^4 - x^5 dx^1$, then (2) reads

$$\begin{split} dD - x^5 dA - E \, dx^1 &= (D_2 - x^5 \, A_2) (dx^2 - x^3 dx^1) + (D_3 - x^5 A_3) (dx^3 - x^4 dx^1) + \\ D_4 (dx^4 - x^5 dx^1) \,, \end{split}$$

or else $E = D_1 - x^5 A_1 + x^3 (D_2 - x^5 A_2) + x^4 (D_3 - x^5 A_3) + x^5 D_4$. When S is 3, that is, when $\omega^3 = dx^1 - x^5 dx^4$, (2) means

$$\begin{split} dA - x^5 dD - E dx^4 &= \lambda (dx^2 - x^3 dx^1) + \mu (dx^3 - x^4 dx^1) + \\ & (A_1 - x^5 D_1 + \lambda x^3 + \mu x^4) (dx^1 - x^5 dx^4) \,, \end{split}$$

where $\lambda = A_2 - x^5 D_2$, $\mu = A_3 - x^5 D_3$. Then $E = x^5 (A_1 - x^5 D_1 + x^3 (A_2 - x^5 D_2) + x^4 (A_3 - x^5 D_3) - D_4)$ (the underlined terms have been forgotten in [G], p. 72⁵).

COROLLARY 1. Keeping a KR pseudo-normal form S fixed, every i. a. Y of the Darboux system $S^{(n-3)}$ has its unique prolongation $[Y]^{n-3} \in L(S)$ s.t. $[[Y]^{n-3}]_{n-3} = Y$. Each i. a. $X \in L(S)$ is obtained in this way: $[[X]_{n-3}]^{n-3} = X$ by the uniqueness of the prolongation.

Therefore, the set L(S) is parametrized – in the way depending on S – by the i.a.'s of Darboux, hence — by smooth functions on \mathbb{R}^3 (contact hamiltonians).⁴

OBSERVATION 1. $L(g^*S) = g_*^{-1}L(S)$ and L(S) is a module over **R**. Hence dim L(S)(0) is an invariant of the classification of germs at $0 \in \mathbf{R}^n$ of KR pseudo-normal forms S on \mathbf{R}^n . We call it the symmetry dimension of the germ of S at 0.

3. Symmetry dimension of members of the exceptional couples of models in dimension 9. Local classification of GC in dimension 9 (obtained by analytic considerations) consists of 93 pairwise non-equivalent models displaying altogether only 89

 $^{{}^{4}}L(S)$ is even a Lie algebra; also the contact hamiltonians can be Poisson-multiplied because of their coming from hamiltonians on the symplectization of $(\mathbf{R}^{3}, \omega^{1})$. Their product is called the *Lagrange bracket*, cf. for inst. [Ly], 1.4.4. In the present paper, however, we do not use those additional structures.

different small growth vectors at $0 \in \mathbf{R}^9$: four pairs are 'served' by growth vectors coinciding in those pairs ([M1], Main Theorem). We mean, sticking to the labels from [M1], the pairs:

(*)1.3.1.2.1 and 1.3.1.2.2, 3.1.3.**2**.1 and 3.1.3.2.2, (**) (***) 3.2.3.3.2 - and 3.2.3.3.2, (****) 3.3.3.2.1 and 3.3.3.2.2.

PROPOSITION 2. (i) Symmetry dimension is an additional invariant that distinguishes geometrically local models in pairs (*), (**), and (****) above.

(ii) Members of (***) (see Ex. 1), however, are not discernible by symmetry dimension.

PROOF. (i). Having recursive formulas for the coordinate functions of i. a.'s of any KR pseudo-normal form (cf. proof of Prop. 1), we iterate those formulas backwards, always evaluating them at 0. In this way, we eventually express i. a.'s at 0 in terms of certain jets at $0 \in \mathbf{R}^3$ of contact hamiltonians. In the particular cases discussed in (i) the outcome is the following, writing horizontally the coordinates of v.f. and denoting, from now on, by |0 the evaluation at 0. For the pair (*):

 $L(1.3.1.2.1) | 0 = \{ [-f_3, f, f_1, f_{11}, f_{111}, 0, 0, -3f_2 - 10f_{13}, 0] | 0, f(x^1, x^2, x^3) \text{ smooth} \};$ $L(1.3.1.2.2) | 0 = \{ [-f_3, f, f_1, f_{11}, f_{111}, 0, 0, -3f_2 - 10f_{13}, -4f_2 - 13f_{13}] | 0, f \text{ smooth} \}.$ In view of the arbitrariness of f (Cor. 1), $L(1.3.1.2.1) | 0 = (\partial_1, \partial_2, \partial_3, \partial_4, \partial_5, \partial_8)$. On the other hand, because $\begin{vmatrix} -3 & -10 \\ -4 & -13 \end{vmatrix} \neq 0$, $L(1.3.1.2.2) \mid 0 = (\partial_1, \partial_2, \partial_3, \partial_4, \partial_5, \partial_8, \partial_9)$. Hence the symmetry dimensions are 6 and 7, respectively.

For the pair (**):

$$L(3.1.3.2.1) | 0 = \{ [-f_3, f, f_1, f_{11}, 0, 0, 0, 5f_2 + 12f_{13}, 0] | 0, f \text{ smooth} \};$$

$$L(3.1.3.2.2) | 0 = \{ [-f_3, f, f_1, f_{11}, 0, 0, 0, 5f_2 + 12f_{13}, 7f_2 + 17f_{13}] | 0, f \text{ smooth} \}.$$

Here dim (L(3.1.3.2.1)|0) = 5, whereas dim (L(3.1.3.2.2)|0) = 6 (because $\begin{vmatrix} 5 & 12 \\ 7 & 17 \end{vmatrix} \neq 0$).

For the pair (****):

$$L(3.3.3.2.1) | 0 = \{ [-f_3, f, f_1, f_{11}, 0, 0, 0, -5f_2 - 13f_{13}, 0] | 0, f \text{ smooth} \};$$

$$L(3.3.3.2.2) | 0 = \{ [-f_3, f, f_1, f_{11}, 0, 0, 0, -5f_2 - 13f_{13}, -7f_2 - 18f_{13}] | 0, f \text{ smooth} \}.$$

So, dim (L(3.3.3.2.1)|0) = 5 and dim (L(3.3.3.2.2)|0) = 6 (as $\begin{vmatrix} -5 & -13 \\ -7 & -18 \end{vmatrix} \neq 0$).

(ii). Proceeding exactly as in (i) one obtains

$$\dim \left(L(\mathfrak{Z}, \mathfrak{Z}, \mathfrak{Z}, \mathfrak{Z}, \mathfrak{Z}, \mathfrak{Z}) \mid 0 \right) = \dim \left(L(\mathfrak{Z}, \mathfrak{Z}, \mathfrak{Z}, \mathfrak{Z}, \mathfrak{Z}) \mid 0 \right) = 6. \blacksquare$$

OPEN QUESTION. How to distinguish geometrically 3.2.3.3.2- from 3.2.3.3.2?

REMARK 1. (a) It is visible that symmetry dimension is a geometric invariant (cf. Obs. 1) just complementary to small growth vector: it takes the same values 5 and 6 on pairs (**) and (****) differing by the small gr. vector. (And not only by it; (**) and (****) differ due to the very elementary reason of Cor. 3.2 of [M1].)

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(b) It is to be underlined that this invariant has already been used for GC in [G]. One of the aims of Gaspar was to distinguish, in a different way than through so-called *reduced tensors* of KR, the five Kumpera–Ruiz local models in dimension 6 ([KR], p. 226–7).

That pioneering use occurred too early, in a sense,⁵ if having already pointed to a limited discernment power of the symmetry dimension alone (taking on three values in that occurrence: 1.1 - 6, 1.3 - 5, 3.1 - 4, 3.2 - 5, 3.3 - 4).

4. The first non-trivial branch of KR pseudo-normal forms in dimension 10. We want to supply some information concerning the true status of KR pseudo-normal forms beyond dimension 9. In dimension 10 and only among codes having just one 3 in the beginning (the simplest non-trivial pseudo-normal KR forms), the situation resembles the general situation with GC in dimension 9: small gr. vector is not the only invariant and the symmetry dimension is only partially complementary to it (cf. Prop. 2). In fact,

THEOREM 3. In dimension 10, here is the list of all different local models among the KR pseudo-normal forms having a unique 3 at the beginning of the code. Small growth vectors at $0 \in \mathbf{R}^{10}$ are given in the second column, with the numbers of repetitions of an integer shown as subscripts. Symmetry dimensions at $0 \in \mathbf{R}^{10}$ are given (only when needed) in the third column.

3.2.1.1.1.1	$[2, 3, 4, 5, 6, 7, 8, 9_2, 10]$	
3.1. 2 .1.1.1	$[2, 3, 4, 5, 6, 7, 8, 9_3, 10]$	
3.1.1. 2 .1. 2 -	$[2, 3, 4, 5, 6, 7, 8, 9_4, 10]$	$\dim(L \mid 0) = 7$
3.1.1. 2 .1.1	$[2, 3, 4, 5, 6, 7, 8, 9_4, 10]$	$\dim(L \mid 0) = 6$
3.1.1.2.1.2	$[2, 3, 4, 5, 6, 7, 8, 9_4, 10]$	$\dim(L \mid 0) = 7$
3.1.1.1. 2 .1	$[2, 3, 4, 5, 6, 7, 8, 9_5, 10]$	
3.1.1.1.1.2	$[2, 3, 4, 5, 6, 7, 8, 9_6, 10]$	
3.1.1.1.1.1	$[2, 3, 4, 5, 6, 7, 8, 9_7, 10]$	

PROOF. That the small growth vector at the origin depends – in the branch in question – only on the place of the first 2 in the code (i.e., on the index of the first non-zero constant in the pseudo-normal form), amounts to a simple calculation. When that first non-vanishing constant is NOT c^8 , its normalization to 1, and reductions to 0 of the constants appearing after it, go in a way similar to [CM] (Thm. 27) and [M1] (Lemmas [32111], [31211]).

For $c^8 \neq 0$, however, it is otherwise. The subsequent constant c^9 can be annihilated by [M1], Lem. [31121], while the last constant c^{10} hides an obstacle. It can only be reduced, when keeping already $c^8 = 1$ and $c^9 = 0$, to exactly one of the values -1, 0, 1 (see the list of proposed local models in Thm. 3). We will start (1°) with the possibility of reduction (simple), then proceed (2°) to the non-equivalence of the three produced values of c^{10} (difficult).

As for the symmetry dimension in these exceptional cases, its computation is much similar to the calculus done in Sec. 3. The results are written in the theorem table – the

⁵Small gr. vector is the only local invariant of GC up to dimension 8 inclusive, as has later been shown in [CM].

symmetry dimension distinguishes 0 from -1 and 1, whereas the last two values are not discernible by this tool.

1°. A reduction of a non-zero $c^{10} = c$ to sgn(c) is straightforward by passing to the bar variables:

$$\begin{aligned} x^{10} &= |c| \,\overline{x^{10}}, \quad x^9 = |c|^{\frac{1}{2}} \,\overline{x^9}, \quad x^4 = |c|^{-\frac{1}{2}} \,\overline{x^4}, \quad x^8 = \overline{x^8}, \quad x^7 = |c|^{-\frac{1}{2}} \,\overline{x^7} \\ x^6 &= |c|^{-1} \,\overline{x^6}, \quad x^5 = |c|^{-\frac{3}{2}} \,\overline{x^5}, \quad x^1 = |c|^{-2} \,\overline{x^1}, \quad x^3 = |c|^{-\frac{5}{2}} \,\overline{x^3}, \quad x^2 = |c|^{-\frac{9}{2}} \,\overline{x^2} \,. \end{aligned}$$

This change of coordinates preserves all but the last Pfaffian equations of the KR writing of 3.1.1.2.1.2, while that last equation assumes the form

$$dx^9 - (c+x^{10})dx^4 = \sqrt{|c|}d\overline{x^9} - \frac{c+|c|\overline{x^{10}}}{\sqrt{|c|}}d\overline{x^4} = \sqrt{|c|}\left(d\overline{x^9} - (\operatorname{sgn}(c) + \overline{x^{10}})d\overline{x^4}\right).$$

 2° . The values -1, 0, and 1 of c^{10} are mutually non-equivalent.

In a rigid situation of keeping $c^8 = 1$ and $c^9 = 0$ fixed, we try to conjugate two a priori arbitrary values c and \tilde{c} of the constant c^{10} .

The distributions D, D dual to the KR pseudo-normal forms in question are spanned by (X, Y_c) and $(X, Y_{\tilde{c}})$, where $X = \partial_{10}$ and, in matrix notation (T – transpose),

$$Y_a^{\mathrm{T}} = \left[x^5, x^3 x^5, x^4 x^5, 1, x^6, x^7, 1 + x^8, x^9, a + x^{10}, 0 \right] \,.$$

For any diffeo $g: (\mathbf{R}^{10}, 0) \hookrightarrow$ sending D to \widetilde{D} , the linear subdistribution (X) common for D and \widetilde{D} is preserved by g (a basic fact, used already in the proof of Prop. 1). Therefore, for certain functions \overline{f} and $\overline{h}, \overline{f} \mid 0 \neq 0$,

$$g_*Y_c = \bar{f}Y_{\tilde{c}} + \bar{h}X \,.$$

Taking this identity at g(x) instead of x gives

$$Dg(x)Y_c(x) = f(x)Y_{\tilde{c}}(g(x)) + h(x)\partial_{10}$$
(3)

for $f = \overline{f} \circ g$, $h = \overline{h} \circ g$; $f \mid 0 \neq 0$. At that, writing $g = (g^1, g^2, \dots, g^{10}), g^1, g^2, g^3$ depend only on x^1, x^2, x^3 , and for $l \geq 4$ g^l depends on x^1, x^2, \dots, x^l (cf. [M1], Cor. 3.1). Moreover,

$$g^{5}(x^{1}, \dots, x^{5}) = x^{5}G(x^{1}, \dots, x^{5})$$
(4)

for some function G, because the hypersurface $E = \{x^5 = 0\}$ can be characterized in invariant terms simultaneously for D and \tilde{D} . In fact, a short calculation shows that both D and \tilde{D} have at points of E the small gr. vectors of the type $[2, 3, 4, 5, 6, 7, 8, 9, 9, \ldots]$ and off E – just the small vector [2, 3, 4, 5, 6, 7, 8, 9, 10]. The preserving of E means $g^5(x^1, x^2, x^3, x^4, 0) = 0$ identically, or else (4).

In the sequel we shall write simply g_k^l for $\frac{\partial g^l}{\partial x^k}$. For instance, the inequality $\frac{\partial g^l}{\partial x^l} | 0 \neq 0$ will henceforth be denoted $g_l^l | 0 \neq 0$.

Equating the coefficients of ∂_9 in (3) (later we just say 'taking scalar equation "9" of (3)') and evaluating the result at 0, a cardinal relation linking c and \tilde{c} reads

$$f^{-1}(g_4^9 + g_7^9 + cg_9^9) | 0 = \tilde{c}.$$

CLAIM. $g_9^9 | 0 = f^{-1} | 0$.

PROOF. In view of (4), eq. "5" of (3) means

$$g^{6} = f^{-1} \left((x^{5})^{2} G_{1} + x^{3} (x^{5})^{2} G_{2} + x^{4} (x^{5})^{2} G_{3} + x^{5} G_{4} + x^{6} (G + x^{5} G_{5}) \right),$$

implying

$$g^6 \in (x^5, x^6) \tag{(\dagger)}$$

 $((x^5, x^6)$ means the *ideal* in \mathcal{F}_0^{10} generated by the functions x^5 and x^6). In turn, eq. "6" of (3) gives g^7 in function of f and g^6 . In fact, by (†),

$$g^7 \equiv f^{-1}(x^6 g_5^6 + x^7 g_6^6) \mod (x^5, x^6),$$

yielding

$$g^7 \in (x^5, x^6, x^7)$$
. (††)

Observe also that g^9 can be expressed by f and g^8 (eq. "8" of (3)), while g^8 gets expressed (eq. "7" of (3)) by g^7 and again f as follows:

$$1 + g^8 = f^{-1} \left(x^5 g_1^7 + x^3 x^5 g_2^7 + x^4 x^5 g_3^7 + g_4^7 + x^6 g_5^7 + x^7 g_6^7 + (1 + x^8) g_7^7 \right) .$$
 (5)

Evaluating this at 0, with (††) implying $g_4^7 \mid 0 = 0$, one gets an important piece of information

$$f^{-1}g_7^7 \mid 0 = 1. (6)$$

As for f, it is explicitly given by eq. "4" of (3) and visibly depends only on x^1, \ldots, x^5 :

$$f = x^5 g_1^4 + x^3 x^5 g_2^4 + x^4 x^5 g_3^4 + g_4^4 \,. \tag{7}$$

Now that we have (5) through (7), we can compute: $g_9^9 | 0 = f^{-1}g_8^8 | 0 = f^{-1} \cdot f^{-1}g_7^7 | 0 = f^{-1} | 0$. Claim is proved.

In view of Claim, the relation linking c and \tilde{c} assumes the form

$$f^{-1}(g_4^9 + g_7^9) + cf^{-2} | 0 = \tilde{c}.$$
(8)

4.1. BASIC LEMMA. $g_4^9 + g_7^9 | 0 = 0$.

PROOF. The fact that $c^8 = 1$ and $c^9 = 0$ are preserved by g means

$$g_4^7 + g_7^7 | 0 = f | 0, \qquad g_4^8 + g_7^8 | 0 = 0.$$
 (a)

On using (5) and (7) for computing $g_7^8 | 0$, and – additionally – the *first* identity in (a) in computing $g_4^8 | 0$, the *second* identity in (a) assumes the form

$$-f_4 + g_6^7 + 2g_{47}^7 | 0 = 0.$$
 (b)

Continuing in this way, writing g^7 in terms of g^6 and f (eq. "6" of (3)), one arrives from (b), after a computation, at

$$-3f_4 + f^{-1}g_5^6 + 3f^{-1}g_{46}^6 | 0 = 0.$$
 (c)

What we eventually need is a relation between f and G. As g^6 can be written in terms of

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f and g^5 (eq. "5" of (3)), another calculation leads from (c) to the (important) identity

$$-3f_4 + 2f^{-2}G_4 | 0 = 0.$$
 (d)

Now, putting $x^5 = x^6 = x^8 = x^9 = 0$ when computing the expression in Basic Lemma (i.e., working with $f^{-1}(g_4^8 + x^7g_6^8 + g_7^8)$ instead of g^9),

$$g_{4}^{9} | 0 = -f_{4}f^{-2}(\underbrace{g_{4}^{8} + g_{7}^{8}}_{= 0 \text{ by (a)}}) + f^{-1}(g_{44}^{8} + g_{47}^{8}) | 0, \quad g_{7}^{9} | 0 = f^{-1}(g_{47}^{8} + g_{6}^{8} + g_{77}^{8}) | 0,$$

so that we have so far

$$g_{4}^{9} + g_{7}^{9} | 0 = f^{-1} (g_{6}^{8} + g_{44}^{8} + 2g_{47}^{8} + g_{77}^{8}) | 0.$$
(9)

Given the summands to-be-computed in (9), we can put $x^5 = x^8 = 0$, i. e., replace g^8 by $-1 + f^{-1}(g_4^7 + x^6g_5^7 + x^7g_6^7 + g_7^7)$. In the course of computations we take into account certain side facts such as $g_{44}^7 \mid 0 = g_{444}^7 \mid 0 = g_{77}^7 = 0$ (caused by (††) and g^7 being affine wrt x^7). As a result, $g_{77}^8 \mid 0 = 2f^{-1}g_{67}^7 \mid 0$, $g_6^8 \mid 0 = f^{-1}(g_{46}^7 + g_5^7 + g_{67}^7) \mid 0$, and, using also the first identity in (a),

$$\begin{split} g^8_{44} &| 0 = (f^{-1})_{44} f + 2(f^{-1})_4 g^7_{47} + f^{-1} g^7_{447} | 0 \\ &= -f^{-1} f_{44} + 2f^{-2} (f_4)^2 - 2f^{-2} f_4 g^7_{47} + f^{-1} g^7_{447} | 0 \,, \\ g^8_{47} &| 0 = -f^{-2} f_4 (g^7_{47} + g^7_6) + f^{-1} (g^7_{447} + g^7_{46}) \,. \end{split}$$

Summing up the above,

$$g_{6}^{8} + g_{44}^{8} + 2g_{47}^{8} + g_{77}^{8} | 0 = f^{-1}(3g_{46}^{7} + g_{5}^{7} + 3g_{67}^{7} - f_{44} + 3g_{447}^{7} - 2f^{-1}f_{4}(-f_{4} + g_{6}^{7} + 2g_{47}^{7})) | 0 = 0.$$

Therefore, by (9) and (b)

$$g_{4}^{9} + g_{7}^{9} | 0 = f^{-2} (3g_{46}^{7} + g_{5}^{7} + 3g_{67}^{7} - f_{44} + 3g_{447}^{7}) | 0.$$
⁽¹⁰⁾

One repeats, basically, the above procedure when reducing the RHS of (10) to g^6 . Firstly, $g_7^7 | 0 = f^{-1}g_6^6 | 0$, or else, by (a) and ($\dagger \dagger$),

$$g_{6}^{6} | 0 = f^{2} | 0.$$
 (e)

In turn,
$$g_5^7 | 0 = f^{-1}g_{45}^6 | 0, g_{67}^7 | 0 = f^{-1}g_{66}^6 | 0 = 0 (g^6 \text{ affine wrt } x^6),$$

 $g_{46}^7 | 0 = -f^{-2}f_4(g_{46}^6 + g_5^6) + f^{-1}(g_{446}^6 + g_{45}^6) | 0,$
 $g_{447}^7 | 0 = (-f^{-2}f_{44} + 2f^{-3}(f_4)^2)g_6^6 - 2f^{-2}f_4g_{46}^6 + f^{-1}g_{446}^6 | 0$
 $\stackrel{\text{by (e)}}{=} -f_{44} + 2f^{-1}(f_4)^2 - 2f^{-2}f_4g_{46}^6 + f^{-1}g_{446}^6 | 0$

Summing up,

$$\begin{aligned} &3g_{46}^7 + g_5^7 + 3g_{67}^7 - f_{44} + 3g_{447}^7 | 0 = \\ & f^{-1}(4g_{45}^6 + 6g_{446}^6 - 4ff_{44} - 3(f_4)^2 - 3f_4(-3f_4 + f^{-1}g_5^6 + 3f^{-1}g_{46}^6)) \,. \end{aligned}$$

Therefore, by (10) and (c),

$$g_{4}^{9} + g_{7}^{9} | 0 = f^{-3} (4g_{45}^{6} + 6g_{446}^{6} - 4ff_{44} - 3(f_{4})^{2}) | 0.$$
⁽¹¹⁾

Clearly, the RHS of (11) can be reduced down to $g^5 = x^5 G$. Firstly $g_6^6 | 0 = f^{-1}G | 0$, or else, by (e),

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$$G \mid 0 = f^3 \mid 0.$$
 (f)

Secondly, $g_{45}^6 | 0 = (f^{-1}G_4)_4 | 0 = -f^{-2}f_4G_4 + f^{-1}G_{44} | 0$. Thirdly,

 $g_{446}^{6} | 0 = (f^{-1}G)_{44} | 0 = (-f^{-2}f_{44} + 2f^{-3}(f_{4})^{2}) G - 2f^{-2}f_{4}G_{4} + f^{-1}G_{44} | 0.$

Taking (f) into account and summing up, $% \left(f^{\prime}\right) =\left(f^{\prime}\right) \left(f^{\prime}\right)$

$$4g_{45}^6 + 6g_{446}^6 - 4ff_{44} - 3(f_4)^2 = f^{-1}(10G_{44} - 10f^2f_{44} + 9f(f_4)^2 - 24f(f_4)^2 - 8ff_4(-3f_4 + 2f^{-2}G_4)) |0.$$

Therefore, by (11) and (d),

$$g_4^9 + g_7^9 |0| = f^{-4} (10G_{44} - 10f^2 f_{44} - 15f(f_4)^2) |0|.$$
(12)

Now G can be eliminated from the RHS of (12). fG is only affine wrt x^4 , so $0 = (fG)_{44} = f_{44}G + 2f_4G_4 + fG_{44}$. This, (d), and (f) yield $G_{44} \mid 0 = -f^2f_{44} - 3f(f_4)^2 \mid 0$. Substituting to (12),

$$g_{4}^{9} + g_{7}^{9} |0 = -5f^{-3}(4ff_{44} + 9(f_{4})^{2}) |0.$$
(13)

The RHS of (13) can be further simplified using (7): the derivatives of f entering (13) can be expressed by those of

$$g^{4} = \frac{fGg^{4}}{fG} = \frac{A(x^{1}, x^{2}, x^{3}) + B(x^{1}, x^{2}, x^{3})x^{4}}{C(x^{1}, x^{2}, x^{3}) + D(x^{1}, x^{2}, x^{3})x^{4}}$$

$$\begin{split} & \text{REMARK 2. i). } A \mid \! 0 = A + B \, x^4 \mid \! 0 = fG \, g^4 \mid \! 0 = 0 \, . \\ & \text{ii). } C \mid \! 0 = fG \mid \! 0 \neq 0 \, . \\ & \text{iii). } f \mid \! 0 = g_4^4 \mid \! 0 = \frac{BC - AD}{C^2} \mid \! 0 = \frac{B}{C} \mid \! 0 \; \text{ by i}) . \\ & \text{iv). } B \mid \! 0 \neq 0 \; \text{by iii).} \\ & \text{v). } f_4 \mid \! 0 = g_{44}^4 \mid \! 0 = -2 \frac{BD}{C^2} \mid \! 0 \, . \\ & \text{vi). } f_{44} \mid \! 0 = g_{444}^4 \mid \! 0 = 6 \frac{BD^2}{C^3} \mid \! 0 \, . \\ & \text{vii). } G \mid \! 0 = \frac{C^2}{B} \mid \! 0 \; (\text{by ii) and iii}) . \end{split}$$

By Rem. 2 ((iii), (v), and (vi)) $4ff_{44} | 0 = 24 \frac{B^2 D^2}{C^4} | 0 = 6(f_4)^2 | 0$ and (13) boils down to

$$g_{4}^{9} + g_{7}^{9} | 0 = -75f^{-3}(f_{4})^{2} | 0.$$
(14)

On the other hand, $G = \frac{fG}{f} = \frac{C + Dx^4}{g_4^4 + x^5(\star)}$ (see (7)). Hence, also by Rem. 2,

$$G_4 \mid 0 = \frac{D}{g4_4} - \frac{Cg_{44}^4}{(g_4^4)^2} \mid 0 = \frac{CD}{B} - C\left(-2\frac{BD}{C^2}\right)\left(\frac{C}{B}\right)^2 \mid 0 = 3\frac{CD}{B} \mid 0$$

Approaching the end, we are going to obtain a formula alternative to (d) for $G_4 | 0$. The identity (f) reads now (see (iii) and (vii) of Rem. 2) $\frac{C^2}{B} | 0 = (\frac{B}{C})^3 | 0$, or else

$$B^4 \mid 0 = C^5 \mid 0.$$
 (g)

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Therefore,

$$3\frac{CD}{B} | 0 = 3\left(\frac{B}{C}\right)^2 \frac{BC^3D}{B^4} | 0 \stackrel{\text{(g)}}{=} 3\left(\frac{B}{C}\right)^2 \frac{BC^3D}{C^5} | 0 = -\frac{3}{2}\left(\frac{B}{C}\right)^2 \left(-2\frac{BD}{C^2}\right) | 0.$$

That is, by (iii) and (v) of Rem. 2,

$$G_4 | 0 = -\frac{3}{2} f^2 f_4 | 0.$$
 (h)

Finally, (d) and (h) together give $f_4 \mid 0 = 0$ and Basic Lemma follows from (14).

4.2. Concluding part of the proof of Thm. 3. In view of Basic Lemma, because of the factor $f^{-2} | 0$ in (8), the values -1, 0, and 1 of c^{10} are mutually non-equivalent as asserted in 2° .

The assessment of classes of local equivalence among the KR pseudo-normal forms 3.1.1.2. (1 or 2).(1 or 2), all covered by the small gr. vector [2, 3, 4, 5, 6, 7, 8, 9₄, 10], is now finished. Thm. 3 is proved.

Addendum. A result obtained after the submission of the present paper underlines clearly a special role of the constant c^{10} in the whole branch of KR pseudo-normal forms 3.1.1.2.1.(1 or 2).(1 or 2)... in dimensions from 11 onwards. (In this notation c^8 is already normalized to 1 and c^9 is annihilated – cf. the beginning of the proof of Thm.3.) Namely, in every dimension n > 10, any such given KR local form is equivalent to exactly one of the following local models: 3.1.1.2.1.2-.1.1...1, 3.1.1.2.1.2.1.1...1,

exactly one of the following local models: $3.1.1.2.1.2-.\underbrace{1.1...1}_{n-10}$, $3.1.1.2\underbrace{1.1...1}_{n-8}$, $3.1.1.2.1.2\underbrace{1.1...1}_{n-10}$. That is, all subsequent constants in this branch are reducible to

0, i.e. are not important. A proof will be included in a paper in preparation (cf. Section 9 of [M2]).

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