

## BLOW-UP ON THE BOUNDARY: A SURVEY

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**Introduction.** In this survey we review some results on blow-up of solutions of the problem

$$(0.1) \quad \frac{\partial u^m}{\partial t} = \Delta u, \quad x \in \Omega, t > 0,$$

$$(0.2) \quad \frac{\partial u}{\partial \nu} = u^p, \quad x \in \partial\Omega, t \geq 0,$$

$$(0.3) \quad u(x, 0) = u_0(x) > 0, \quad x \in \bar{\Omega},$$

$$(0.4) \quad \frac{\partial u_0}{\partial \nu} = u_0^p, \quad x \in \partial\Omega,$$

where  $m, p > 0$  and  $\Omega$  is either a smoothly bounded domain in  $\mathbb{R}^N$  or  $\Omega = \mathbb{R}_+^N = \{(x_1, x') : x' \in \mathbb{R}^{N-1}, x_1 > 0\}$ ,  $\nu$  is the outward normal.

Over the past two decades this problem has received considerable interest. For  $\Omega$  bounded,  $m = 1$  and  $p > 1$  it was shown by Levine and Payne ([LP1]) in 1974 and by Walter ([Wa]) in 1975 that there are solutions which blow up in finite time. This means that

$$\limsup_{t \rightarrow T} \max_{\bar{\Omega}} u(x, t) = \infty \quad \text{for some } T < \infty.$$

The major questions that have been studied since then are:

1. *For which values of  $m, p$  does blow-up occur?*
2. *For which initial functions does blow-up occur?*
3. *Where are the blow-up points located?*
4. *With which rate (in  $t$ ) does the solution approach the blow-up time?*

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5. *What is the profile (in  $x$ ) at the blow-up time?*
6. *Can blow-up in infinite time occur?*
7. *If  $\Omega = \mathbb{R}_+^N$ , what is the critical Fujita exponent?*

Here we give a survey of answers (or partial answers) to the above questions and we also present some basic ideas under the simplest circumstances.

1.  $m = 1$ ,  $\Omega = (0, \infty)$ . We will assume throughout this section that  $u_0 \in C^1$  and  $\lim_{x \rightarrow \infty} u_0(x) = 0$ .

1.1. *If  $p \leq 1$  then all solutions are global.*

To see this it is sufficient to verify that

$$v(x, t) = e^{\alpha^2 t} (e^{-\alpha x} + c)$$

is a supersolution if  $c > 0$  and  $\alpha^2 = (1 + c)^p$  (cf. [GL, Remark 2.2]). Let us mention here that if  $p < 1$  then uniqueness fails to hold (cf. [DFL, Theorem 3.5]).

1.2. *If  $p > 1$  then there are explicit selfsimilar solutions that blow up in finite time.*

They are of the form

$$u(x, t) = (T - t)^{-\lambda} f_-(\xi), \quad T > 0, \quad \lambda = \frac{1}{2(p-1)}, \quad \xi = \frac{x}{\sqrt{T-t}},$$

$f_-$  is the unique bounded solution of

$$\begin{aligned} f_-''(\xi) - \frac{\xi}{2} f_-'(\xi) - \lambda f_-(\xi) &= 0, \quad \xi > 0, \\ -f_-'(0) &= f_-^p(0). \end{aligned}$$

The function  $f_-$  is given explicitly in terms of degenerate hypergeometric functions (see [FQ, Lemma 3.1]) and it is not difficult to verify that  $u$  has the following properties (cf. [DFL, Lemma 3.1]):

- (i)  $u_t > 0$  in  $(0, \infty) \times (0, T)$ ,
- (ii)  $u(x, T) = kx^{-2\lambda}$ ,  $k = \pi^{-1/2} \left( \lambda \frac{\Gamma^p(\lambda + 1/2)}{\Gamma(\lambda + 1)} \right)^{2\lambda}$ ,
- (iii)  $x^{2\lambda} u(x, t) \rightarrow k$  as  $x \rightarrow \infty$ ,  $0 \leq t \leq T$ .

Here  $x = 0$  is the only blow-up point and the blow-up rates in  $t$  and  $x$  are  $(T - t)^{-\lambda}$  and  $x^{-2\lambda}$ , respectively. We shall show that many solutions behave similarly.

1.3. *If  $u$  is a solution that blows up in the time  $T$  and*

$$\beta = \inf_{x>0} -\frac{u_0'(x)}{u_0^p(x)} \in (0, 1]$$

then

$$\limsup_{t \rightarrow T} u(x, t) \leq [\beta(p-1)]^{-2\lambda} x^{-2\lambda} \quad \text{for } x > 0.$$

This was shown in [B, Theorem 2].

The proof follows by a simple maximum principle argument. If we take

$$J(x, t) = u_x + \beta u^p$$

then it is not difficult to show that  $J \leq 0$  in  $\Omega \times (0, T)$ . If we integrate the inequality

$$u_x + \beta u^p \leq 0$$

we obtain the assertion.

1.4. *Assume  $u$  blows up at the time  $T$  and  $u'_0 \leq 0$ . Then there is a  $\delta = \delta(u_0) > 0$  such that*

$$\limsup_{t \rightarrow T} u(x, t) \geq p^{-\frac{p}{p-1}} (p-1) x^{-2\lambda} \quad \text{for } x \in (0, \delta).$$

To show this one uses the intersection-comparison method as in [GKS]. Namely, for any  $u_0$  there are  $\alpha_0, \delta > 0$  such that the stationary solution  $U_\alpha(x) = -\alpha^p x + \alpha$  has in  $(0, \delta)$  a unique intersection with  $u_0$  for all  $\alpha \geq \alpha_0$  and  $U_\alpha(0) > u_0(0)$ . Since  $u_x \leq 0$  and  $u$  blows up, we obtain that for any  $\alpha \geq \alpha_0$  there is a  $t_\alpha \in (0, T)$  such that  $U_\alpha(0) < u(0, t_\alpha)$ . The number of intersections is nonincreasing therefore it is actually equal to zero at  $t = t_\alpha$ . Hence  $\limsup_{t \rightarrow T} u(x, t) \geq \sup_{\alpha \geq \alpha_0} U_\alpha(x)$  for  $x \in (0, \delta)$  and it is easy to verify that  $\sup_{\alpha \geq \alpha_0} U_\alpha(x) = p^{-p/(p-1)} (p-1) x^{-2\lambda}$ .

In 1.3 and 1.4 we described the profile in  $x$  and next we turn to the same question but in  $t$ .

1.5. *Assume  $u_0 \in C^3$ ,  $(-1)^i u_0^{(i)} \geq 0$ ,  $i = 1, 2, 3$  and  $-u_0'''(0) = pu_0^{p-1}(0)u_0''(0)$ . Then  $u$  blows up at a finite time  $T$  and*

$$u(0, t) \leq (p-1)^{-\lambda} (T-t)^{-\lambda} \quad \text{for } t \in (0, T).$$

We proceed as in [FQ, Lemma 2.1] (cf. also [DFL, Theorem 3.4]). By the maximum principle  $u, u_t \geq 0$  and  $u_x, u_{xt} \leq 0$ . Using this and integration by parts we obtain

$$\begin{aligned} \frac{1}{2} u^{2p}(0, t) &= \frac{1}{2} u_x^2(0, t) = - \int_0^\infty u_{xx}(x, t) u(x, t) dx \\ &= - \int_0^\infty u_t(x, t) u_x(x, t) dx \\ &= - \lim_{x \rightarrow \infty} u_t(x, t) u(x, t) + u_t(0, t) u(0, t) + \int_0^\infty u_{xt}(x, t) u(x, t) dx \\ &\leq u_t(0, t) u(0, t). \end{aligned}$$

From the inequality

$$u_t(0, t) \geq \frac{1}{2} u^{2p-1}(0, t)$$

we conclude that  $u$  blows up at a time  $T$  and integrating over  $(t, T)$  we obtain the result.

1.6. *Assume that  $u'_0 \leq -u_0^p$  and  $u$  blows up at a finite time  $T$ . Then*

$$u(0, t) \geq (\lambda p^{-1})^\lambda (T-t)^{-\lambda} \quad \text{for } t \in (0, T).$$

We proceed similarly as in 1.3. By the maximum principle,  $J(x, t) = u_x + u^p \leq 0$  in  $\Omega \times (0, T)$  and  $J(0, t) = 0$ . Therefore  $J_x(0, t) = u_t(0, t) - pu^{2p-1}(0, t) \leq 0$ . Integration of the last inequality over  $(t, T)$  yields the result. Notice that 1.5 and 1.6 give upper and

lower bounds for  $T$  in terms of  $p$  and  $u_0(0)$ . As an example of a function  $u_0$  satisfying all assumption in 1.5 and 1.6 we can take

$$u_0(x) = [4\lambda a(x+a)^{-2}]^{2\lambda}, \quad a > 0.$$

For the existence time  $T$  of the solution starting from this initial function we obtain

$$\frac{(p-1)a^2}{8p} \leq T \leq \frac{(p-1)a^2}{4}.$$

1.7. *If  $p > 2$  then there are global selfsimilar solutions. They are of the form*

$$u(x, t) = (t_0 + t)^{-\lambda} f_+(\zeta), \quad \zeta = \frac{x}{\sqrt{t_0 + t}}, \quad t_0 > 0,$$

$f_+$  satisfies

$$\begin{aligned} f_+''(\zeta) + \frac{\zeta}{2} f_+'(\zeta) + \lambda f_+(\zeta) &= 0, \quad \zeta > 0, \\ -f_+'(0) &= f_+^p(0), \end{aligned}$$

and it can be expressed explicitly in terms of degenerate hypergeometric functions (cf. [DFL]).

1.8. *If  $p \in (1, 2]$  then all solutions blow up in finite time. If  $p > 2$  then there are both global and nonglobal solutions. ( $p = 2$  is the critical Fujita exponent.)*

The first statement is shown by Kaplan type arguments in [GL]. The second one follows from 1.7.

1.9. *Assume  $p > 2$ . Then the solution blows up in finite time provided*

$$\liminf_{x \rightarrow \infty} x^{2\lambda} u_0(x) \geq k,$$

$k$  is from 1.2(ii). On the other hand, there are global solutions such that  $\lim_{x \rightarrow \infty} x^{2\lambda} u(x, t)$  exists and is positive for all  $t > 0$ .

The first assertion follows by comparison with selfsimilar solutions from 1.2. The property from the second statement is satisfied for a one parameter family of selfsimilar solutions from 1.7 (cf. [DFL]).

1.10. *If  $\Omega = \mathbb{R}_+^N$  then the Fujita type result from 1.8 holds with the critical exponent  $p = 2$  replaced by  $p = 1 + 1/N$  (cf. [DFL]).*

**2.**  $m < 1$ ,  $\Omega = (0, \infty)$ . Assume  $\sup |(u_0^{m^{-1}-1})'| < \infty$ ,  $u_0$  has compact support and  $-u_0'(0) = u_0^p(0)$ .

2.1. *If  $p \leq (m+1)/2$  then all solutions are global. If  $p > (m+1)/2$  then there are solutions that blow up in finite time.*

2.2. *If  $p \in ((m+1)/2, m+1]$  then all solutions blow up in finite time. If  $p > m+1$  then global solutions exist.*

All statements in 2.1 and 2.2 were proved in [GL]. The most difficult and very interesting result here is blow-up of all solutions when  $p = m+1$ . All other results in 2.1 and 2.2 are proved by comparison with sub- and supersolutions of selfsimilar type.

**3.  $m = 1$ ,  $\Omega$  is a bounded domain in  $\mathbb{R}^N$**

3.1. *If  $p \leq 1$  then all solutions are global. If  $p > 1$  then there are solutions that blow up in finite time.*

Blow-up of solutions emanating from “large” initial data was established in [LP1] using energy methods. In [Wa] both the global existence and the blow-up result were shown by comparison arguments.

3.2. *If  $p > 1$  then all (positive) solutions blow up in finite time.*

We indicate here how this fact follows from the result in [Fa] (discussed below) which says that global solutions are bounded provided  $p < N/(N - 2)$  if  $N > 2$ . It is easy to see that there are no positive steady states and that zero is unstable. If a solution were global then it would be bounded and its  $\omega$ -limit set would have to contain nonnegative steady states — a contradiction. If  $p \geq N/(N - 2)$  then a comparison argument finishes the proof. In [LMW1] this result was established for balls in  $\mathbb{R}^N$  and simply connected domains in  $\mathbb{R}^2$ . See also [HY] for a short proof.

3.3. *If  $a \in \overline{\Omega}$  is a blow-up point then  $a \in \partial\Omega$ . (We call  $a$  a blow-up point if there are  $\{x_n\} \subset \Omega$  and  $t_n \rightarrow T < \infty$  such that  $x_n \rightarrow a$  and  $\lim_{t \rightarrow T} u(x_n, t) = \infty$ .)*

This result was first proved for radially symmetric solutions in [LMW1] using a maximum principle argument similar as in 1.3. The general case was settled later in [HY] under the assumption that  $u \leq C(T - t)^{-q}$  for some  $C, q > 0$ . (This is satisfied for example if  $\Delta u_0 \geq 0$ .)

3.4. *There is an example of single point blow-up on the boundary.*

This example can be found in [H2].

3.5. *Assume that  $\partial\Omega \in C^{2+\alpha}$  and  $p < N/(N - 2)$  if  $N > 2$ . Suppose  $u_0 \in C^2(\overline{\Omega})$  and  $\Delta u_0 \geq 0$  in  $\Omega$ . Then*

$$\max_{\overline{\Omega}} u(x, t) \leq C(T - t)^{-\lambda},$$

$\lambda = 1/2(p - 1)$  as in Section 1.

This result was first established in the radially symmetric case (no restriction on  $p$  is needed there) in [FQ] under additional assumptions on  $u_0$  (cf. 1.5). In [HY] the general case was proved under a stronger restriction on  $p$ , namely,  $p < (N - 1)/(N - 2)$  if  $N > 2$ . This restriction was needed because of lack of a sharp nonexistence result for

$$\begin{aligned} \Delta u &= 0 && \text{in } \mathbb{R}_+^N, \\ \frac{\partial u}{\partial x_1} &= u^p && \text{for } x_1 = 0. \end{aligned}$$

The sharp nonexistence result was established later in [H1].

3.6. *Suppose  $\partial\Omega \in C^{1+\alpha}$ . Then*

$$\max_{\overline{\Omega}} u(x, t) \geq c(T - t)^{-\lambda}.$$

Using an integral representation of  $u$ , this was shown in [HY].

3.7. Assume  $\Omega = (-1, 1)$ ,  $u_0(x) = u_0(-x)$  and  $u_0^{(i)}(x) \geq 0$ ,  $i = 1, 2, 3, 4$ ,  $x \in [0, 1]$ . Let  $T$  be the blow-up time. Then for any  $y \geq 0$  we have

$$(T - t)^\lambda u(1 - y\sqrt{T - t}, t) \rightarrow f_-(y) \quad \text{as } t \rightarrow T$$

uniformly on compact intervals;  $f_-$  is from 1.2.

For the proof (also in the radial case on balls in higher dimension) we refer to [FQ]. For a generalization see [HY].

3.8. Suppose that  $\partial\Omega \in C^{2+\alpha}$  and

$$\max_{\overline{\Omega}} u(x, t) \leq C(T - t)^{-\lambda}$$

for some  $C > 0$ . If for some  $K > 0$

$$\liminf_{t \rightarrow T} (T - t)^\lambda \inf_{|y| \leq K} u(a + y\sqrt{T - t}, t) = 0,$$

then  $a$  is not a blow-up point.

This nondegeneracy of the blow-up limit was established in [H2].

3.9. Let  $u$  be a global solution of

$$\begin{aligned} u_t &= \Delta u, & x \in \Omega, t > 0, \\ \frac{\partial u}{\partial \nu} &= f(u), & x \in \partial\Omega, t > 0, \\ u(x, 0) &= u_0(x), & x \in \overline{\Omega}, \end{aligned}$$

with  $\partial\Omega \in C^2$  and  $f \in C^\alpha$  for some  $\alpha \in (0, 1)$ . Suppose

$$uf(u) \geq (2 + \varepsilon) \int_0^u f(v) dv - C$$

for some positive constants  $\varepsilon, C$ . Assume

- (i)  $|f(u)| \leq g(u)$  for some increasing  $C^1$  function  $g$  if  $N = 1$ ,
- (ii)  $|f(u)| \leq g(\vartheta)e^{\vartheta u^2}$  for some positive function  $g$  and all  $\vartheta > 0, u \in \mathbb{R}$  if  $N = 2$ ,
- (iii)  $|f(u)| \leq \vartheta|u|^{N/(N-2)} + g(\vartheta)$  for some positive function  $g$  and all  $\vartheta > 0, u \in \mathbb{R}$  if  $N > 2$ .

Then  $u$  is uniformly bounded in  $C^{1,\alpha}$ .

This was proved in [L]. It is a significant improvement of the result from [Fa]. It says that under the above assumptions there are just two possible types of behavior of solutions:

- (a) blow-up in finite time,
- (b) global existence and uniform boundedness.

Blow-up in infinite time cannot occur.

**4.  $m > 0$ ,  $\Omega$  is a bounded domain in  $\mathbb{R}^N$**

4.1. Assume  $(N - 2)m < N + 2$ . If  $\limsup_{t \rightarrow T} \max_{\overline{\Omega}} u(x, t) = \infty$  then also

$$\limsup_{t \rightarrow T} \int_{\partial\Omega} |u(X, t)|^r dS = \infty \quad \forall r > (N - 1)(p - 1).$$

It was proved in [Fo1] (for a more general reaction term and with no sign restriction on  $u_0$ ) that for any  $r > (N - 1)(p - 1)$  there exist positive constants  $M, \xi$ , independent of  $T$  such that

$$|u(x, t)| \leq M(1 + \sup_{\overline{\Omega}} |u_0(x)|) \left( 1 + \sup_{0 \leq \tau \leq t} \int_{\partial\Omega} |u(X, \tau)|^r dS \right)^\xi$$

$\forall (x, t) \in \overline{\Omega} \times [0, T)$  and the assertion follows.

The proof of the above estimate is based on Moser's iteration technique and it makes use of the inequalities:

$$\int_{\partial\Omega} |u|^{p+\lambda} \leq \left( \int_{\partial\Omega} |u|^{\frac{(\lambda+1)(N-1)}{N-2}} \right)^P \left( \int_{\partial\Omega} |u|^{(N-1)(p-1)+\varepsilon} \right)^Q \left( \int_{\partial\Omega} |u|^{\lambda+1} \right)^R,$$

$$P = \frac{(p-1)(N-2)}{(N-1)(p-1)+\varepsilon}, \quad Q = \frac{p-1}{(N-1)(p-1)+\varepsilon}, \quad R = \frac{\varepsilon}{(N-1)(p-1)+\varepsilon}$$

if  $N > 2$ , and

$$\int_{\partial\Omega} |u|^{p+\lambda} \leq \left( \int_{\partial\Omega} |u|^{2(\lambda+1)(p-1+\varepsilon)/\varepsilon} \right)^P \left( \int_{\partial\Omega} |u|^{p-1+\varepsilon} \right)^Q \left( \int_{\partial\Omega} |u|^{\lambda+1} \right)^R,$$

$$P = \frac{\varepsilon(p-1)}{(p-1+\varepsilon)(2(p-1)+\varepsilon)}, \quad Q = \frac{p-1}{p-1+\varepsilon}, \quad R = \frac{\varepsilon}{2(p-1)+\varepsilon}$$

if  $N = 2$  ( $0 < \varepsilon < \infty$ ).

4.2. Assume  $\Omega = (-1, 1)$ ,  $u_0(x) = u_0(-x)$  and

(i)  $0 < m < 1$ .

If  $p \leq m$  then each solution exists globally whereas in the case  $m < p$  all solutions blow up in finite time. In the case  $m < p \leq 1$  solutions become unbounded on the whole space interval  $[-1, 1]$ , but for  $p > 1$  the only blow up points are  $x = \pm 1$ .

(ii)  $m \geq 1$ .

If  $2p \leq m + 1$  then all solutions are global and for  $2p > m + 1$  all solutions blow up in finite time.

All statements in 4.2 except for the case  $2p = m + 1 > 2$  were proved in [Fo2]. The borderline case  $2p = m + 1 > 2$  was settled later in [Wo] (see 4.3 below). The results are proved by comparison with solutions emanating from special chosen initial data (cf. [Fo2]). In some cases also the rate in  $t$  and profile in  $x$  at the blow-up time are shown.

If  $0 < m < p \leq 1$  and  $u$  is a solution that blows up in the time  $T$  such that  $u_x, u_{xx}$

are nonnegative on  $[0, 1]$ ,  $u_0(1) > 1$  then

$$\xi^p \leq T(p-m)m^{-1} \left( \int_0^1 u_0^m(x) dx \right)^{(p-m)/m} \leq 1,$$

$\xi = 1 - u_0^{p-1}(1)$  and

$$\xi^{(2p-m)/(p-m)} \leq c(T-t)^{1/(p-m)} u(x, t) \leq \xi^{-1}$$

$\forall (x, t) \in [0, 1] \times [0, T]$ ,  $c = ((p-m)/m)^{1/(p-1)}$ .

If  $0 < m < 1 < p$  and  $u$  is a solution such that  $u_x$  is nonnegative on  $[0, 1]$  then

$$u(x, t) \leq \frac{C}{(1-|x|)^{1/(p-1)}}$$

for some positive constant  $C$ .

If  $2p > m + 1 > 2$  and  $u$  is a solution that blows up at time  $T$  such that  $u_x, u_{xx}, u_{xxx}$  are nonnegative on  $[0, 1] \times [0, T]$  then

$$\frac{C_\varepsilon}{(T-t)^{1/(2p-m-1+\varepsilon)}} \leq u(1, t) \leq \frac{C}{(T-t)^{1/(2p-m-1)}}$$

for some positive constants  $C, C_\varepsilon$  and  $0 < \varepsilon \ll 1$ .

The results of [Fo2] were generalized by [Wo] in two ways. In [Wo] general nonlinearities are allowed and the domain is an  $N$ -dimensional ball or any simply connected smooth domain in  $\mathbb{R}^2$ .

#### 4.3. The problem

$$\begin{aligned} u_t &= \Delta \Phi(u) && \text{in } B_R \times (0, T), \quad B_R = \{x \in \mathbb{R}^N : |x| < R\}, \\ \frac{\partial \Phi(u)}{\partial \nu} &= f(u) && \text{on } S_R \times [0, T], \quad S_R = \{x \in \mathbb{R}^N : |x| = R\}, \\ u(x, 0) &= u_0(x) > 0 && \text{in } B_R, \end{aligned}$$

where  $\Phi, f$  are increasing functions that are positive for  $u$  positive together with their derivatives and which go to infinity as  $u$  goes to infinity, was studied in [Wo]. It was shown that

- (A) if  $\Phi'(u) \geq C > 0$  and
  - (i)  $f(u)/(1+u)$  is bounded then all solutions are global,
  - (ii)  $\int^\infty ds/f(s) < \infty$  then every solution blows up in finite time,
- (B) if  $0 < \Phi'(u) \leq C$  and
  - (i)  $\Phi$  is concave or  $f(u)/\Phi(u)$  is nondecreasing and  $\sqrt{\Phi'(u)}f(u)/\Phi(u)$  is bounded then every solution exists globally,
  - (ii)  $\Phi$  is concave,  $\liminf_{u \rightarrow \infty} f(u)\sqrt{\Phi'(u)}/\Phi(u) > 0$  and

$$\int^\infty \frac{\sqrt{\Phi'(s)} ds}{f(s)} < \infty$$

then each solution blows up in finite time.



### 5. Related problems

#### 5.1. The problem

$$(5.1) \quad u_t = \nabla(a(u)\nabla u), \quad x \in \Omega, t > 0,$$

$$(5.2) \quad \frac{\partial u}{\partial \nu} = 1, \quad x \in \partial\Omega, t > 0,$$

$$(5.3) \quad u(x, 0) = u_0(x) > 0, \quad x \in \bar{\Omega},$$

was studied in [Y], where  $a \in C^1$  is such that  $a, a' > 0$  and  $\limsup_{u \rightarrow \infty} a'(u)/a(u) < \infty$ . It was shown in [Y] that all solutions are global if and only if  $\int^\infty ds/a(s) = \infty$ . Also, some results on the profile near blow-up were established.

If we take  $a(u) = m^{-1}u^{\frac{1}{m}-1}$ ,  $0 < m \leq 1/2$  and  $v = u^{1/m}$  then  $v$  satisfies

$$(5.4) \quad (v^m)_t = \Delta v, \quad x \in \Omega, t > 0,$$

$$(5.5) \quad \frac{\partial v}{\partial \nu} = \frac{1}{m}v^{1-m}, \quad x \in \partial\Omega, t > 0,$$

which is a special case of (0.1), (0.2) (if we neglect the factor  $1/m$  in (5.5)).

5.2. In [WW], the boundary condition (5.2) was replaced by

$$\frac{\partial u}{\partial \nu} = b(u)$$

and a global existence – global nonexistence result was proved.

5.3. In [LP2], the Laplace operator in (0.1) was replaced by an elliptic operator of order  $2k$ , and (0.2) was changed to correspond to the elliptic operator. For that problem with  $m = 1$ , a “large” data blow-up result was established.

5.4. In [LS], the homogeneous Dirichlet condition was prescribed on a part of the boundary and “large” data blow-up was shown for  $m = 1$ .

5.5. In [CFQ], [LMW2] and [Q] the following problem with a damping term in the equation was considered:

$$u_t = \Delta u - au^p, \quad x \in \Omega \subset \mathbb{R}^N, t > 0,$$

$$\frac{\partial u}{\partial \nu} = u^q, \quad x \in \partial\Omega, t > 0,$$

$$u(x, 0) = u_0(x) \geq 0, \quad x \in \bar{\Omega},$$

with  $\Omega$  bounded,  $p, q > 1$  and  $a > 0$ . For  $N = 1$  it was shown in [CFQ] that

(i) if  $p < 2q - 1$  (or  $p = 2q - 1$  and  $a < q$ ) then there are solutions which blow up in finite time,

(ii) if  $p > 2q - 1$  (or  $p = 2q - 1$  and  $a > q$ ) then all solutions are global and bounded,

(iii) if  $p = 2q - 1$  and  $a = q$  then all nontrivial solutions exist globally but they are not bounded, they tend (as  $t \rightarrow \infty$ ) pointwise to a singular steady state.

The statements (i) and (ii) were proved in [CFQ] also for balls in higher dimension. But for a general domain  $\Omega$  only some partial results can be found in [CFQ]. It was shown later in [Q] for a general domain  $\Omega$  that

(a) if  $p < 2q - 1$  (or  $p = 2q - 1$  and  $a$  is small) then there are solutions which blow up in finite or infinite time,

(b) if  $p > 2q - 1$  then all solutions are global and bounded.

5.6. In [DFL] the following system was studied:

$$\begin{aligned} u_t &= \Delta u, & v_t &= \Delta v, & x &\in \mathbb{R}_+^N, t > 0, \\ -\frac{\partial u}{\partial x_1} &= v^p, & -\frac{\partial v}{\partial x_1} &= u^p, & x_1 &= 0, t > 0, \\ u(x, 0) &= u_0(x) \geq 0, & v(x, 0) &= v_0(x) \geq 0, & x &\in \mathbb{R}_+^N, \end{aligned}$$

with  $p, q > 0$ . It was shown there (among other things) that blow-up may occur if and only if  $pq > 1$  and all nontrivial solutions blow up if and only if

$$\max\left(\frac{p+1}{pq-1}, \frac{q+1}{pq-1}\right) \geq N.$$

When we referred to [DFL] in Section 1, we did that with the hope that interested readers will easily see how to modify the results (or proofs) in the easier scalar case.

5.7. In [FL] the authors studied the profile of solutions that quench on the boundary. They studied the problem

$$\begin{aligned} u_t &= u_{xx}, & x &\in (0, 1), t > 0, \\ u_x(0, t) &= 0, & t &> 0, \\ u_x(1, t) &= -u^{-\beta}(1, t), & x &\in [0, 1], \end{aligned}$$

with  $\beta > 0$ . Every solution of this problem reaches zero (quenches) in finite time.

5.8. The heat equation with a condition similar to (0.2) prescribed on a hypersurface  $\Gamma$  in a bounded domain  $\Omega$  was studied in [CY]. Sufficient condition for global existence and finite time blow-up were established there and also some results on the blow-up rate and blow-up set were proved.

5.9. Assume  $0 < m, r < \infty$ . The problem

$$\begin{aligned} (|u|^{m-1}u)_t &= \sum_{i=1}^N (|u_{x_i}|^{r-1}u_{x_i})_{x_i} & x &\in \Omega, t > 0, \\ \nabla_r u \cdot \nu &= f(u) & x &\in \partial\Omega, t > 0, \\ u(x, 0) &= u_0(x), & u_0 &\in L^\infty(\Omega) \cup W_{r+1}^1(\Omega), \end{aligned}$$

where  $\nabla_r u = (|u_{x_1}|^{r-1}u_{x_1}, \dots, |u_{x_N}|^{r-1}u_{x_N})$  and

$$f(u)\text{sign } u \leq L(|u|^p + 1), \quad L \geq 0, \quad 0 \leq p < \infty,$$

was studied in [Fo3]. It was shown that if

$$q > \max\left\{1, \frac{N-1}{r}\right\} \max\{p-r, 0\}$$

then there exists a positive function  $\mathcal{F} \in C^2(\mathbb{R}_+^2)$  depending solely on the data and  $q$

such that

$$\|u(t)\|_{L^\infty(\Omega)} \leq \mathcal{F}\left(\|u_0\|_{L^\infty(\Omega)}, \operatorname{ess\,sup}_{0 \leq \tau \leq t} \int_{\partial\Omega} |u(X, \tau)|^q dS\right)$$

for a.e.  $t \in [0, T]$  ( $\mathcal{F}(x, y) \rightarrow \infty$  if  $y \rightarrow \infty$ ).

The global existence result was proved under the following assumptions:

$$p \leq \min\{m, r\} \quad \text{or} \quad r < p < \frac{r(m+1)}{r+1}$$

and  $p < p^*$ , where  $p^* = r(m+2)$  if  $N = 1$  and  $p^* = r(N + \max\{p, m\} + 1)/N$  if  $N \geq 2$ .

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